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**UMR CNRS 6143
Morphodynamique Continentale et Côtière
M2C
Université de Caen, Basse Normandie**

1. Effect of cold on the rocks and soils
Mechanical effects of freeze-thaw cycles



2. Influence of rheological parameters on the mechanical properties of geomaterials.
Movement of materials and the process of erosion.



3. Interaction of surface waves and turbulence with bottom and suspended particles



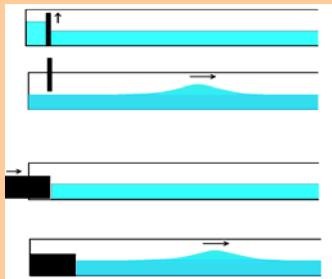
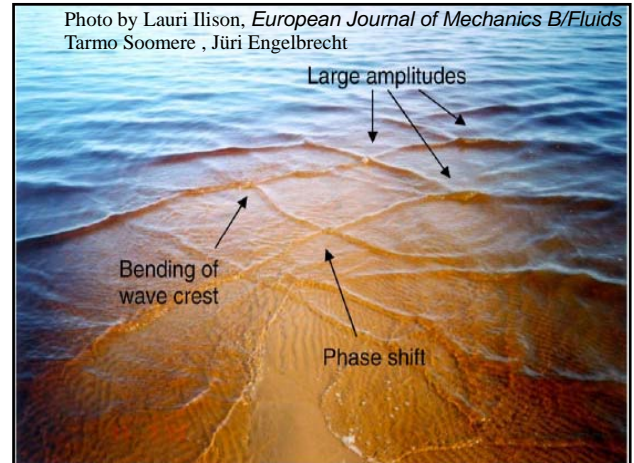
Several wave flumes equipped with wave wakers

**Interaction of solitons with sandy
bottom in a shallow water resonator**

Outline

1. Introduction.
2. Excitation of solitons in a shallow water resonator.
 - 2.1 Map of regimes
 - 2.2 Numerical simulation of soliton using Boussinesq equation
3. Interaction of solitons with sandy bottom.
 - 3.1 Evolution of sandy bottom profile
 - 3.2 Theoretical description
4. Segregation of particles under the action of solitons.
5. Conclusions.

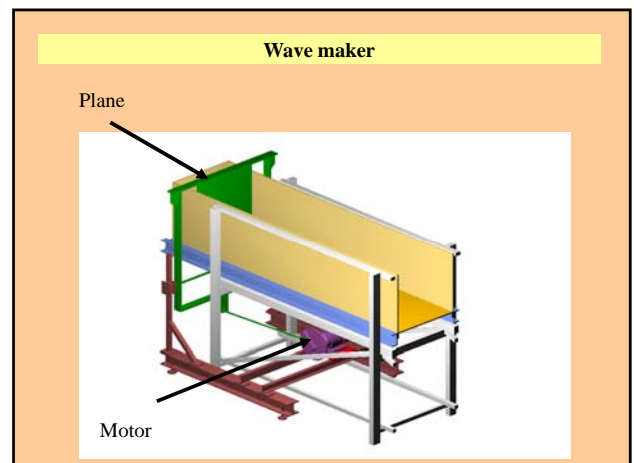
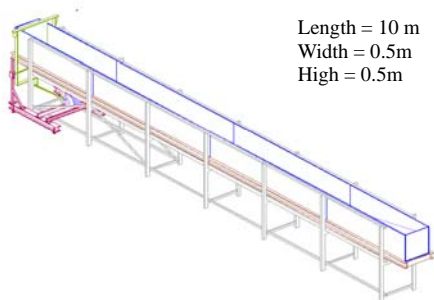
1. Introduction.

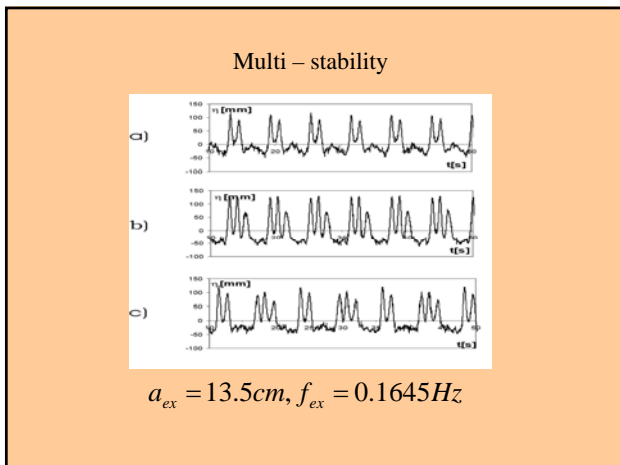
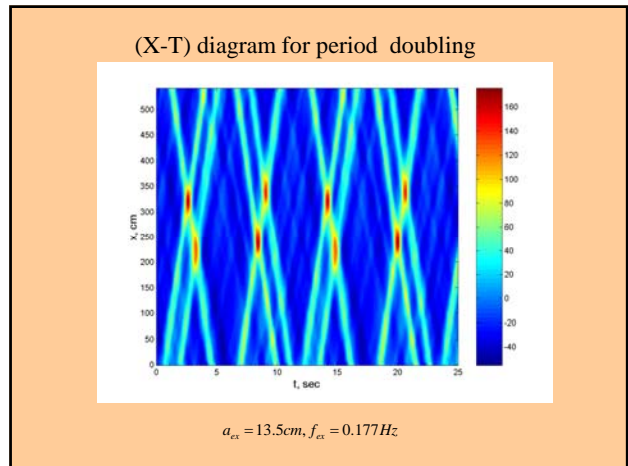
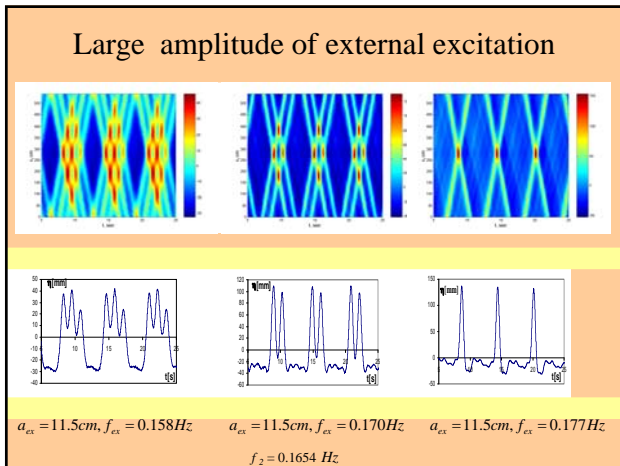
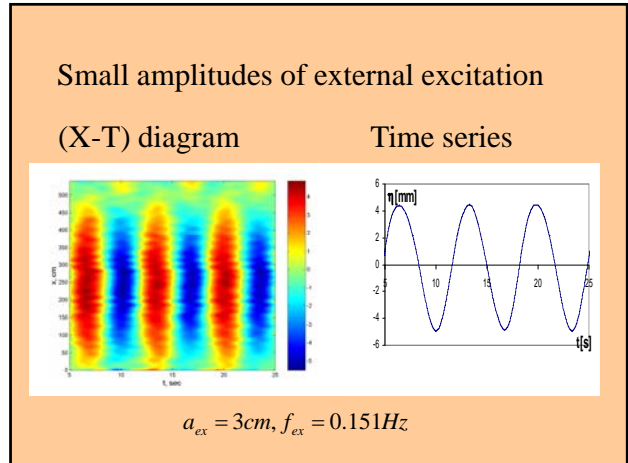
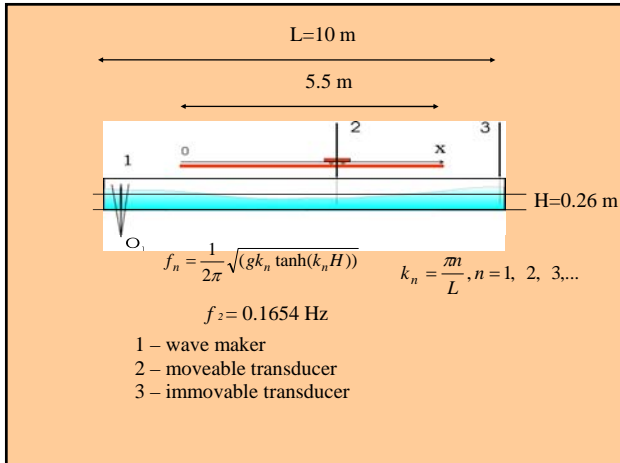


Hammack&Segur 1974
Olsen, Smith&Scott 1984
Kuwabara, Hasegawa&
Kono 1986

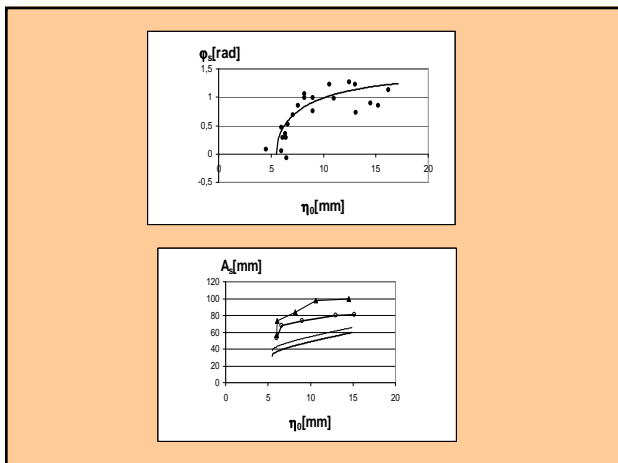
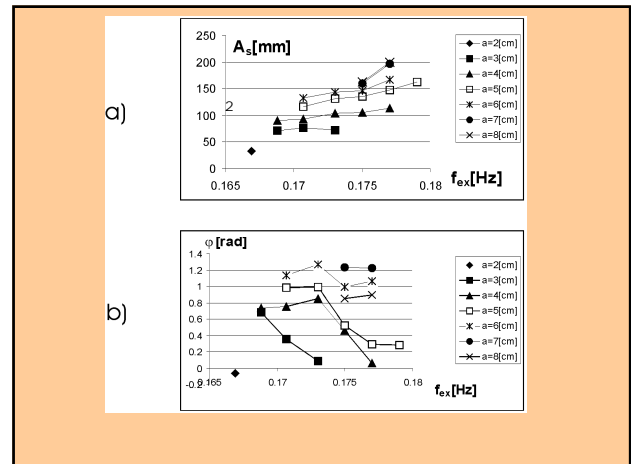
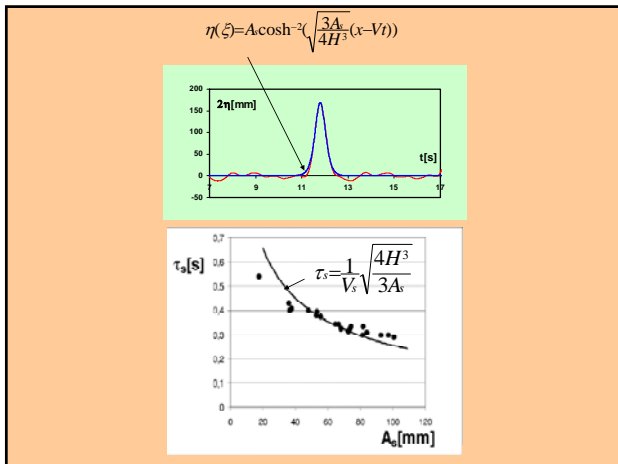
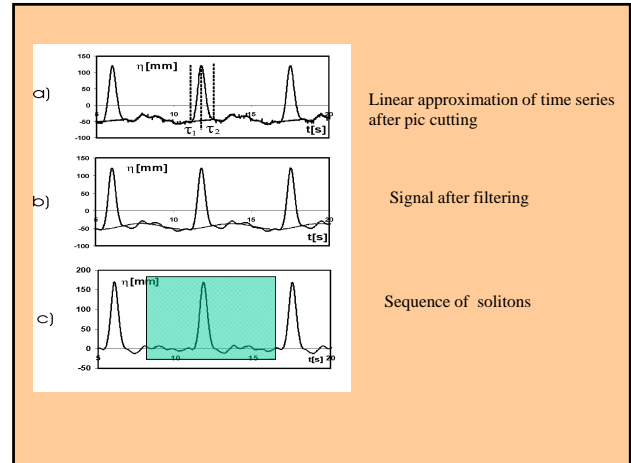
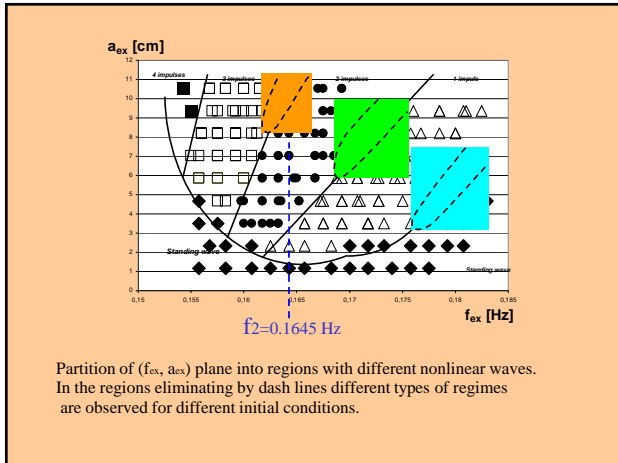
Maxworthy 1976
Weidman&Maxworthy 1978

2. Excitation of solitons in a shallow water resonator.





2.1 Map of regimes

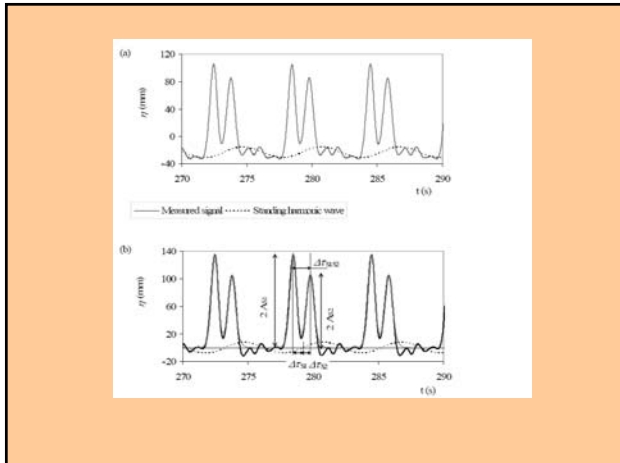


- Gorshkov, Ostrovsky and Papko (1979)
- Ostrovsky and Potapov (1999): Studies for solitons of electromagnetic waves in resonators (section of an LC line excited by sinusoidal voltage)
- Analogy with the solitons of electromagnetic waves → Equations for the amplitude and phase of the soliton

$$\frac{dE_s}{dt} = \frac{3\omega\eta_0}{2H} E_s \cos\phi_s - \delta E_s$$

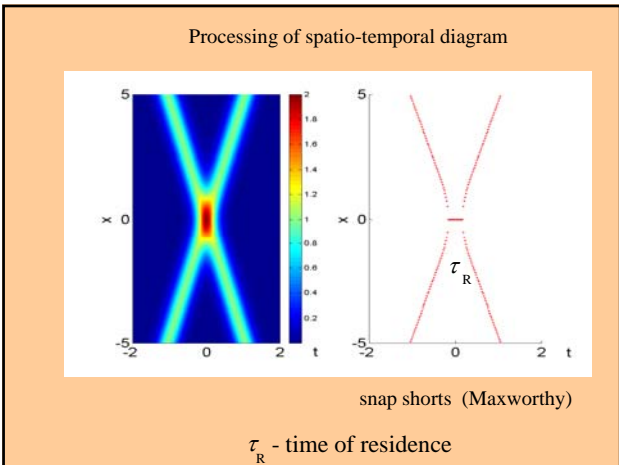
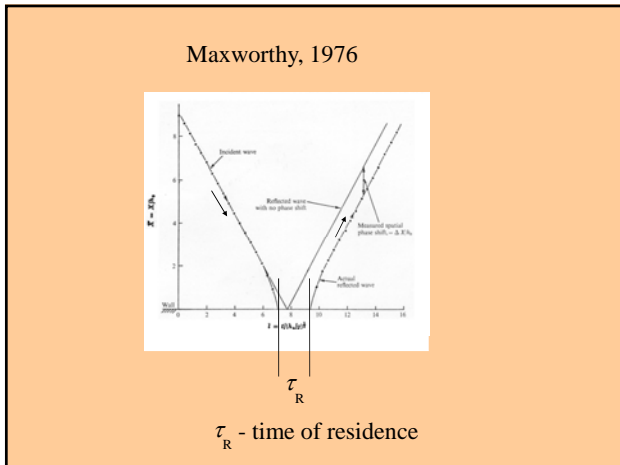
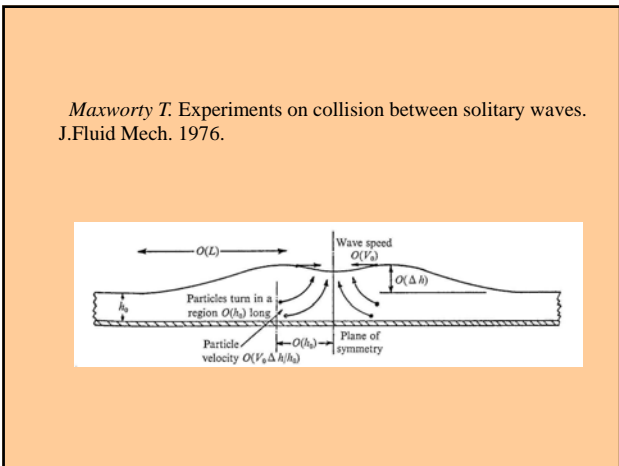
$$\frac{d\phi_s}{dt} = \frac{\sqrt{gH}}{2} \frac{A_s k}{H} - \frac{3}{2} \eta_0 k \sqrt{\frac{g}{d}} \sin\phi_s - \Delta$$

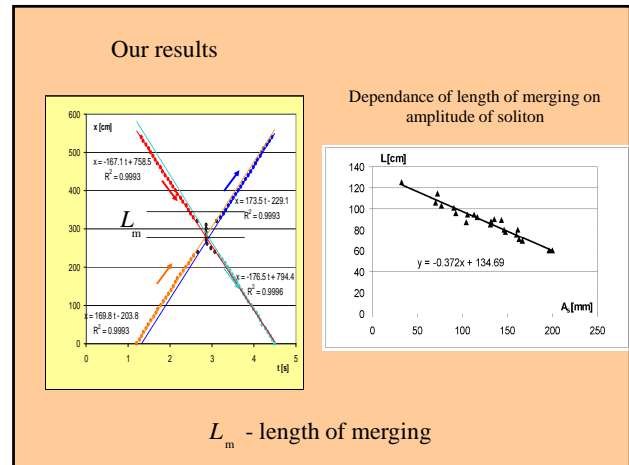
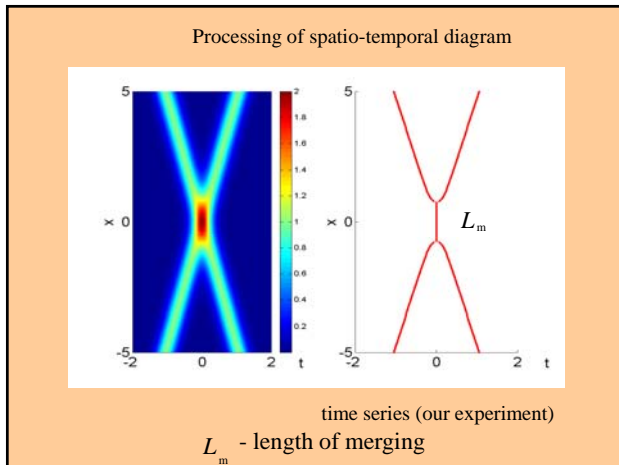
$E_s = \int_{-\infty}^{+\infty} \eta_s^2 dx \sim A_s^{3/2}$ δ : dissipation of the soliton



Test Number	f (Hz)	h_0 (cm)	A_{S1} (mm)	A_{S2} (mm)	ϕ_{S1} (rad)	ϕ_{S2} (rad)	A_N (mm)	Number of pulses
1	0.172	7.5	70.9	-	0.82	-	5.1	1
2	0.173	6.5	75.3	-	0.91	-	5.3	1
3	0.173	6	70.3	-	1.07	-	5.2	1
4	0.165	6	60.2	43.6	0.66	-0.66	3.6	2
5	0.167	6	61.7	48.5	0.70	-0.70	4.3	2
6	0.173	6	81.8	-	1.09	-	4.5	1
7	0.167	6	67.8	52.5	0.80	-0.57	4.0	2
8	0.167	7	89.0	70.7	0.81	-0.53	4.2	2

Comparison with results of Maxworthy





2.2 Numerical simulation of soliton using Boussinesq equation.

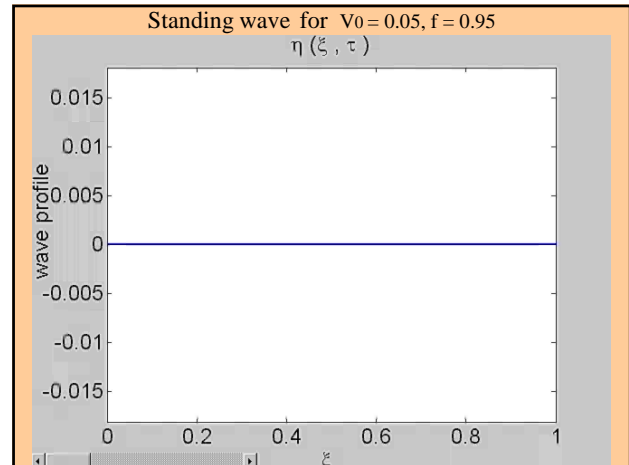
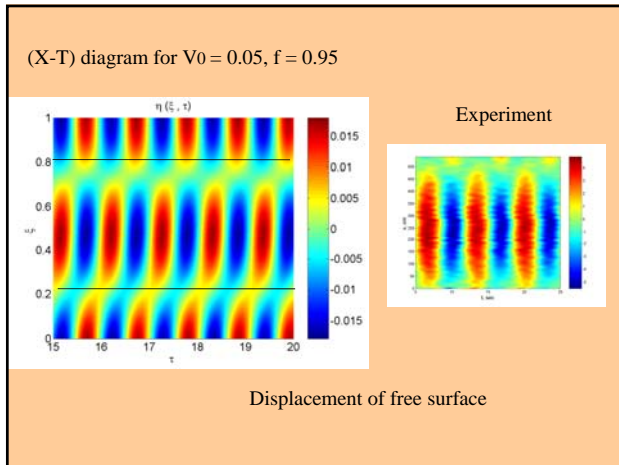
Boussinesq approximation for nonlinear waves of free surface in shallow water (η : displacement of surface, v : horizontal velocity)

$$\eta_\tau + v_\xi + \varepsilon(v\eta)_\xi = \frac{1}{6}\mu^2 v_{\xi\xi\xi} + \gamma_1 \eta_{\xi\xi} - \gamma\eta$$

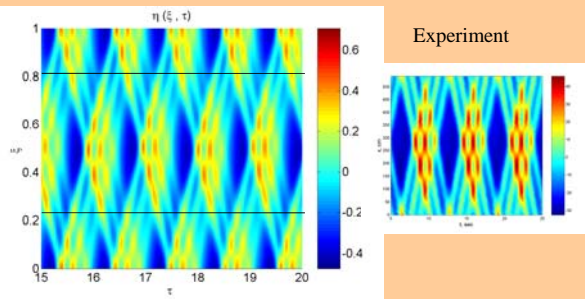
$$v_\tau + \eta_\xi + \varepsilon v v_\xi = \frac{1}{2}\mu^2 v_{\xi\xi\xi} + \gamma_1 v_{\xi\xi} - \gamma v$$

Initial conditions: (i) $\eta(\tau=0, \xi) = 0$ et $v(\tau=0, \xi) = 0$
 (ii) bruit

Boundary conditions: $\frac{\partial \eta}{\partial \xi}(\tau, \xi=0) = 0$
 $v(\tau, \xi=0) = V_0 \sin(2\pi f\tau)$
 $\frac{\partial \eta}{\partial \xi}(\tau, \xi=1) = 0$ $v(\tau, \xi=1) = 0$

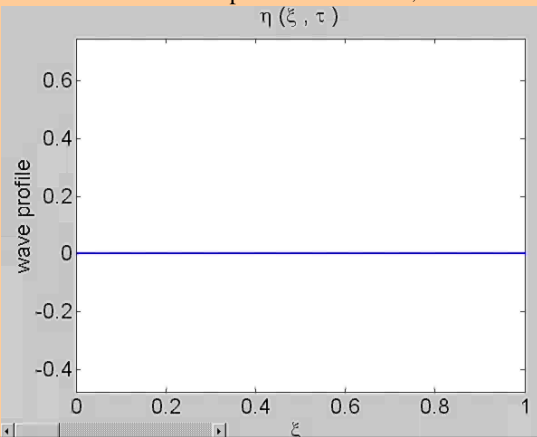


(X-T) diagram for $V_0=0.5, f=0.95$
(three impulses in one period)

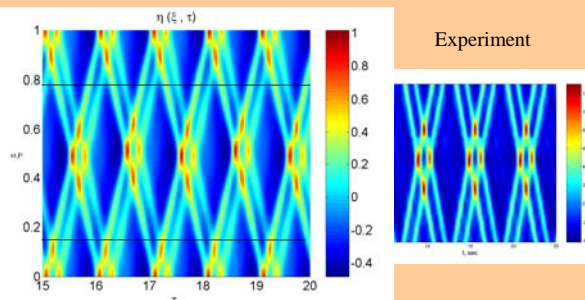


Displacement of free surface

3 solitons on period for $V_0=0.5, f=0.95$

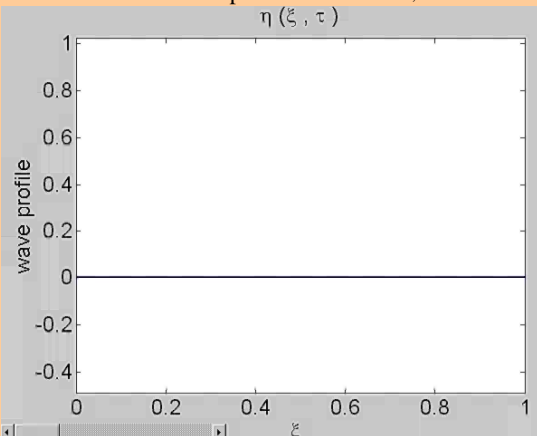


(X-T) diagram for $V_0=0.5, f=0.98$
(2 solitons in one period)

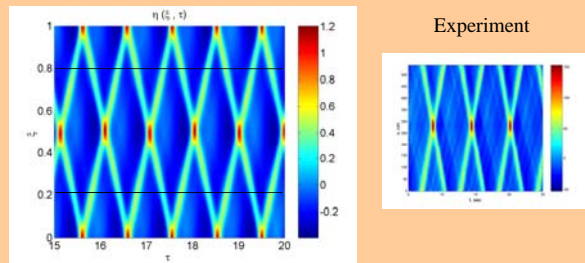


Displacement of free surface

2 solitons on period for $V_0=0.5, f=0.98$

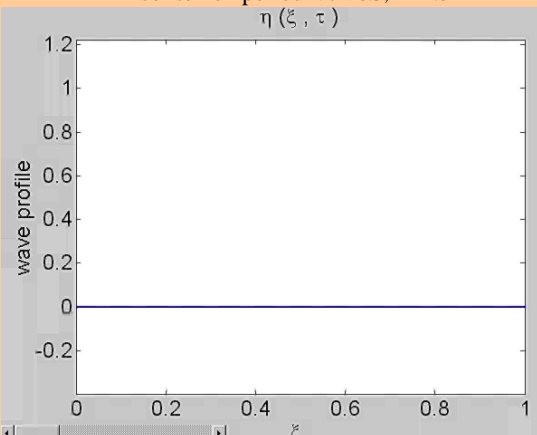


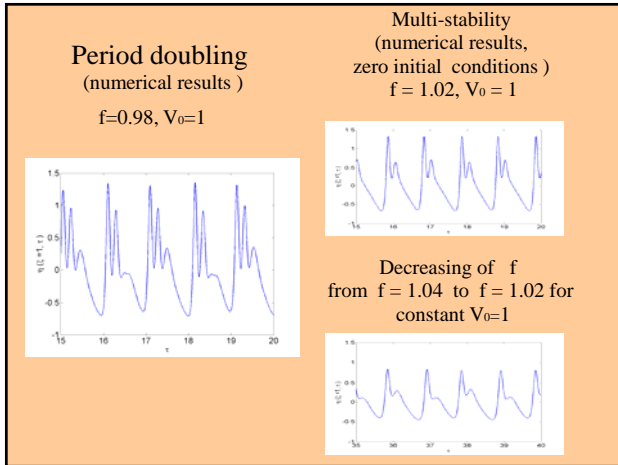
(X-T) diagram for $V_0=0.5, f=1.02$
(one soliton for period)



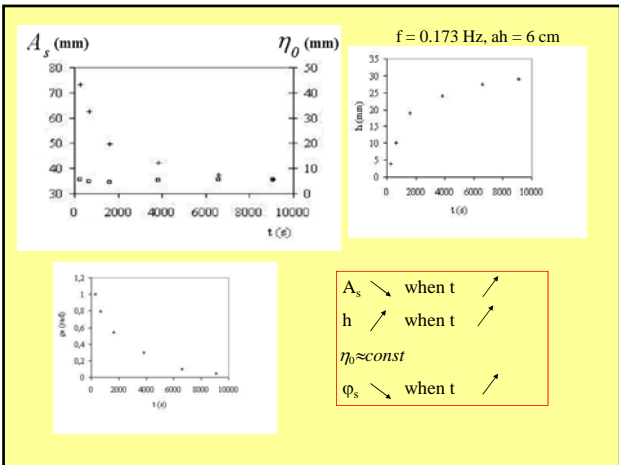
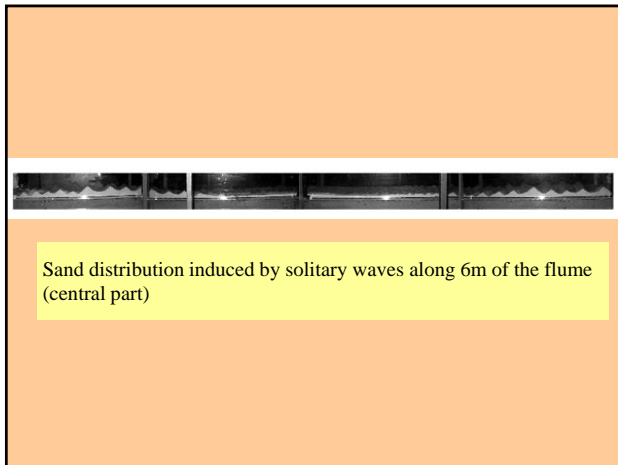
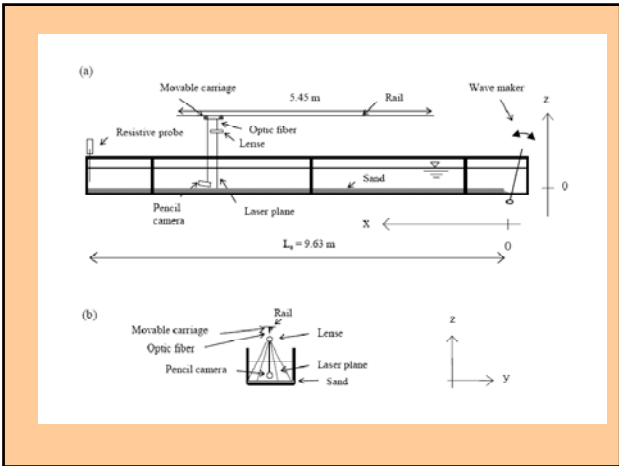
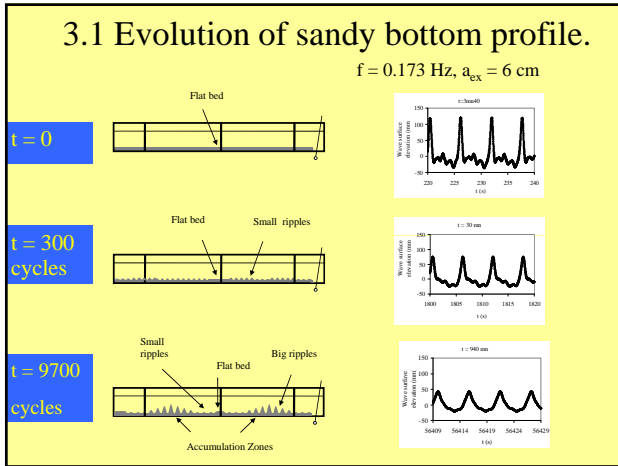
Displacement of free surface

1 soliton on period $V_0=0.5, f=1.02$





3. Interaction of solitons with a sandy bottom.



Ripple wavelength: $L \sim 10$ cm
 Harmonic wave wavelength: $L_h = 9.63$ m
 Negligible scattering of the harmonic wave by the ripples as $L \ll L_h$ ($L/L_h \sim 10^{-2}$)
 $\Rightarrow \eta_0$ is not affected by the ripples

Significant decrease of the amplitude of soliton by:

- organized vortices and turbulence (dissipation)
- solitary wave scattering

Equations for the amplitude and phase of the soliton

$$\frac{dE_s}{dt} = \frac{3\omega\eta_0}{2H} E_s \cos\phi_s - (\delta + ah)E_s$$

$$\frac{d\phi_s}{dt} = \frac{\sqrt{gH}}{2} \frac{Ak}{H} \frac{3}{2}\eta_0 k \sqrt{\frac{g}{d}} \sin\phi_s - \Delta$$

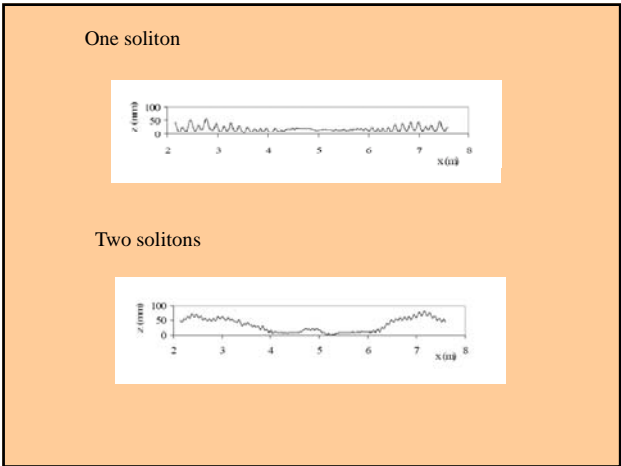
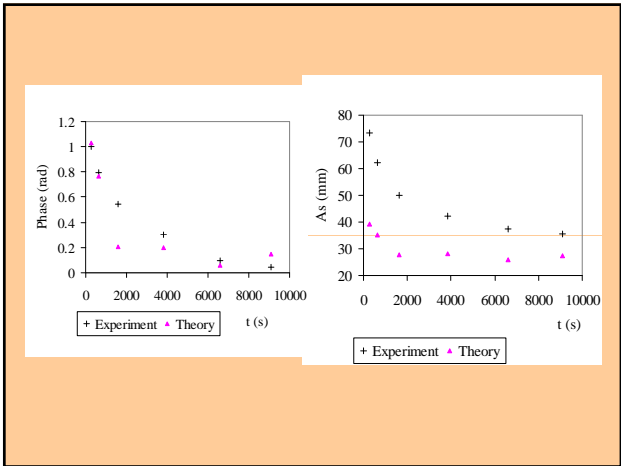
α : phenomenological coefficient for the dissipation of the soliton due to sand ripples

steady state:

$$\phi_s = \arccos\left[\frac{2H(\delta + ah)}{3\omega\eta_0}\right]$$

$$A_s = \frac{2\Delta}{k} \sqrt{\frac{H}{g}} + 3\eta_0 \sqrt{1 - \left(\frac{2H(\delta + ah)}{3\omega\eta_0}\right)^2}$$

Dissipation of the solitary waves due to sand ripples
 $\Rightarrow \phi_s \searrow, A_s \searrow$



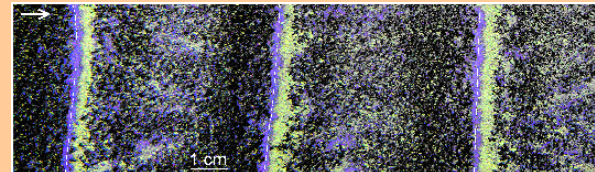
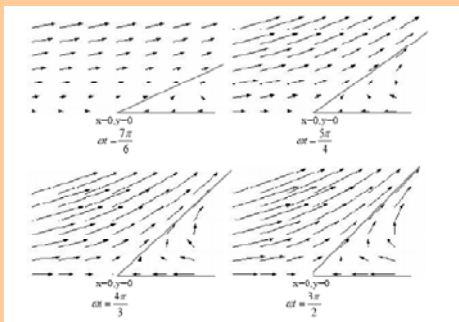
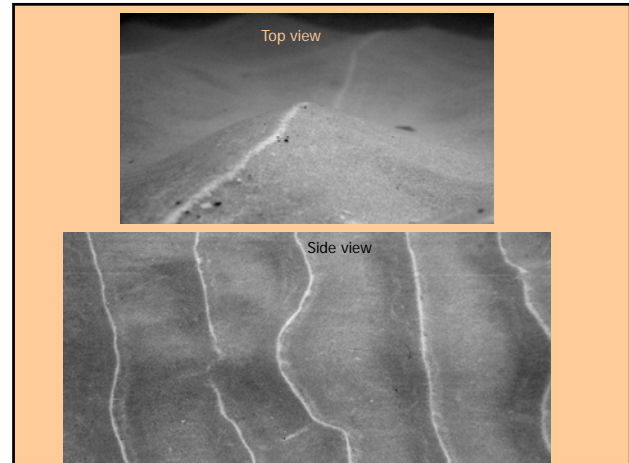
Bar formation

The figure shows a graph of surface elevation z (mm) versus x (m) for standing waves, with two peaks labeled 'Node' and two troughs labeled 'Antinode'. Below the graph is a schematic diagram of bar formation showing the relationship between surface elevation and bed topography.

Solitary waves \rightarrow Bars crests are located beneath the nodes of surface elevation

Standing waves \rightarrow Bar crests are located beneath the antinodes of surface elevation

4. Segregation of particles under the action of solitons.



5. Conclusions

- Solitons and bound state of solitons were revealed in a wave flume used in resonant mode.
- Strong interaction between sandy bottom and non-linear surface waves occurs.
- Segregation of sinking particles in oscillating flow was found

Bibliography

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- F. Marin, A. B. Ezersky Formation dynamics of sand bedforms under solitons and bound states of solitons in a wave flume used in resonant mode". *Europ. J. Mech.*, 2007, v.27, p.251-267.
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